## Novel Heavy Flavor Suppression Mechanisms in the Quark-gluon Plasma

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uantum chromodynamics, the fundamental theory of strong interactions, predicts that at exceedingly high temperatures and densities ordinary nuclear matter should melt down in a plasma-like state of quarks and gluons. One of the important experimental signatures of the quark-gluon plasma (QGP) creation in heavy ion collisions is the detailed suppression pattern of high transverse momentum ( $p_T$ ) hadrons, dubbed jet quenching [1]. Jet quenching for light mesons, such as  $\pi$ , K, and  $\eta$ , at the Relativistic Heavy Ion Collider (RHIC) is well explained by radiative energy loss calculations. One such example of neutral pion suppression in Au+Au and Cu+Cu reactions is shown in Fig.1. Our theoretical calculations of the QGP-induced suppression  $R_{AA}$  versus  $p_T$  and the colliding system size show good agreement with the PHENIX experimental data.

In contrast, recent data on the suppression of single non-photonic electrons, coming from the decay of heavy charm (D) and beauty (B) mesons, cannot be reproduced by medium-induced gluon bremsstrahlung phenomenology in conjunction with a physically reasonable set of QGP temperatures and densities. Recently, we identified a possible reason for this striking discrepancy [2]. The model for predicting the quenched hadron spectra in A+A collisions, where finalstate-in-medium effects are important, has been very simple: it is assumed that the hard jet hadronizes in vacuum, having fully traversed the region of dense nuclear matter,  $L_{OCP}^{T} \sim 6$  fm. It is easily demonstrated that this assumption is not valid for the heavy charm and beauty quarks and mesons. In fact, for  $p_{\tau} = 10 \text{ GeV}$ they tend to form at times t = 0.4 fm and t = 1.5 fm, respectively, well within the hot and dense QGP phase. Consequently, the natural mechanism of heavy flavor suppression is a competition between D- and B-meson dissociation and c- and b-quark fragmentation, both of which emulate energy loss. Our predictions for the suppression  $R_{AA}(p_{\tau})$  of  $e^++e^-$  coming from the semi-leptonic heavy flavor decays are compared in Fig. 2 with the STAR and PHENIX data [3]. Thus, our approach presents an elegant resolution to the "heavy flavor puzzle" at RHIC.

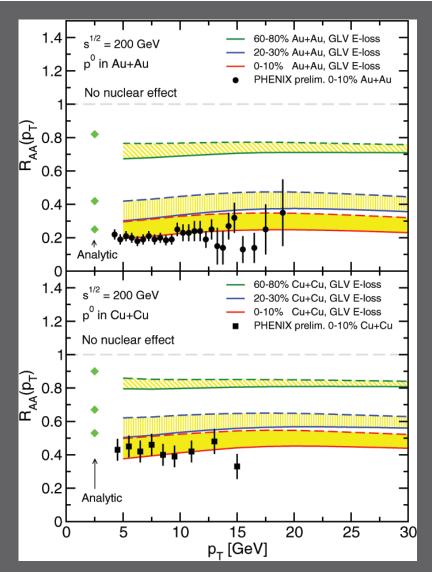


Fig. 1. Top panel: neutral pion quenching predictions versus  $p_{\tau}$  and centrality in Au+Au collisions at RHIC compared with the PHENIX central suppression data at C.M. energy of 200 GeV. Bottom panel: similar comparison is done in Cu+Cu reactions.

A natural *unique* consequence [4] of this theory is that B-mesons are attenuated as much as D-mesons at transverse momenta as low as  $p_{\rm T}\sim 10$  GeV for the conditions prevalent in central Au+Au collisions at C.M. energy of 200 GeV (Fig. 3). Our model will be put to a critical test by experimental advances aiming at direct and separate measurements of D- and B-meson  $R_{AA}(p_{\rm T})$  [5].

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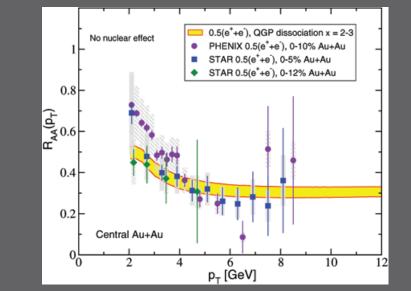


Fig. 2. Suppression of inclusive non-photonic electrons from D- and B-meson spectra softened by collisional dissociation in central Au+Au collisions. Data on non-photonic electron quenching is from PHENIX and STAR.

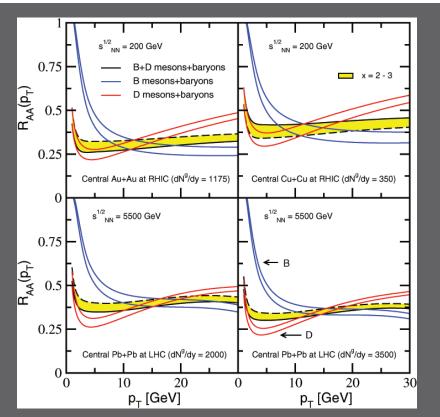


Fig. 3. Individual suppression of D- and B-meson production via collisional dissociation in the QGP. Top panels show numerical results for  $R_{AA}(p_{\tau})$  in central Au+Au and Cu+Cu collisions at RHIC. Bottom panel shows predictions for central Pb+Pb collisions at the LHC.